

# An Improved Self-Powered Switching Interface for Piezoelectric Energy Harvesting

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**Abstract**—In piezoelectric energy harvesting (PEH), with the use of the nonlinear technique named *synchronized switching harvesting on inductor* (SSHI), the harvesting efficiency can be greatly enhanced. Furthermore, the introduction of its self-powered feature makes this technique more applicable for stand-alone systems. In this article, a modified circuitry and an improved analysis for self-powered SSHI are proposed. With the modified circuitry, direct peak detection and better isolation among different units within the circuit can be achieved, both of which result in further removal on dissipative components. In the improved analysis, details in open circuit voltage, switching phase lag, and voltage inversion factor are discussed, all of which lead to a better understanding to the working principle of the self-powered SSHI. Both analyses and experiments show that, in terms of harvesting power, the higher the excitation level, the closer between self-powered and ideal SSHI; at the same time, the more beneficial the adoption of self-powered SSHI treatment in piezoelectric energy harvesting, compared to the standard energy harvesting (SEH) technique.

## I. INTRODUCTION

Structured by a certain number of spatially distributed autonomous sensor devices, the wireless sensor networks (WSNs) have been widely used in military, industrial and civilian applications for monitoring different physical and environmental conditions. Many concerns on energy consumption were addressed due to their nature on wireless connections [1]. Different energy harvesting techniques have been investigated for the purpose to broaden energy sources, alleviate the dependence on batteries, and hopefully someday make all WSN devices self-powered [2].

Piezoelectric energy harvesting (PEH) is one of the promising techniques that can scavenge energy from ambient vibration sources. With its electromechanical coupling characteristic, a piezoelectric element can generate electricity when strain is produced. Since the deformation in vibrating structures is alternating, the generated electricity is also alternating. An interface circuitry is needed for AC-DC conversion. The standard energy harvesting (SEH) interface involves only a bridge rectifier for the AC-DC conversion [3], [4].

To further enhance the energy harvesting capability, Guyomar et al. proposed a nonlinear treatment named *synchronized switching harvesting on inductor* (SSHI) [5]. It was claimed that, under the same displacement excitation, SSHI can increase the harvested power by several hundred percents, compared to SEH [5], [6]. To implement SSHI, a

displacement sensor and a controller were usually needed for synchronization and generation of switching commands [5]–[7], until the emergence of its self-powered version. Based on their experiment, Lallart and Guyomar claimed that the self-powered SSHI can harvest 1.6 times more of power than SEH [8]. They have also considered the influence of voltage gaps, which are produced by diodes and transistors in the circuitry, over the harvested power. Yet, the influences of other components, e.g. the capacitance of the envelope detector, have not been pointed out. Besides, two important parameters, i.e. the switching delay phase  $\varphi$  and inversion factor  $\gamma$  were regarded as constants [8]. But in fact, these two parameters are constants only when the open circuit voltage  $V_{OC}$  (related to maximum displacement) and storage voltage  $V_{DC}$  (voltage across  $C_{rect}$ ) are constants. A more complete analysis should take these into account, in order to compare the energy harvesting efficiencies under different  $V_{OC}$  and  $V_{DC}$ .

In this article, the working principles of piezoelectric energy harvesting with SSHI interface are introduced in Section II. A modified circuit topology for self-powered SSHI is proposed in Section III. With the adoption of complementary topology, the peak detection unit can be separated from the switching inductive shunt; therefore, better isolation can be achieved. In Section IV, referring to the analysis in [8], an improved analysis to the self-powered SSHI is proposed. It provides a more comprehensive understanding to the self-powered SSHI. Experiments are conducted; experimental and analytical results are compared and discussed in Section V. At the end, a conclusion is given in Section VI.

## II. PRINCIPLES

### A. Piezoelectric Energy Harvesting

Given a typical piezoelectric device, e.g. a piezoelectric cantilever with shunt circuit, its schematic representation is shown in Fig. 1(a). Under harmonic displacement excitation, the piezoelectric element can be modeled as a equivalent current source  $i_{eq}$  in parallel with the piezoelectric clamped capacitance  $C_p$ , as shown in Fig. 1(b).  $i_{eq}$  is proportional to the velocity  $\dot{x}$  with the relation of

$$i_{eq}(t) = \alpha_e \dot{x}(t), \quad (1)$$

where  $\alpha_e$  is the piezoelectric force-voltage coupling factor.

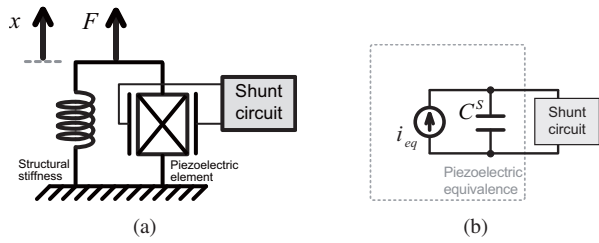


Fig. 1. Schematic representation and equivalent circuit of a typical piezoelectric device. (a) Schematic representation. (b) Equivalent circuit.

Different shunt circuits for different treatments are designed to extract energy from the vibrating mechanical structure. For the purpose of structural damping, the extracted energy is directly dissipated; while for energy harvesting, a portion of the extracted energy is reclaimed and stored in electrical form for subsequent usage.

### B. SSHI

In a vibrating piezoelectric device, since the induced voltage across the piezoelectric capacitance is alternating, the most conventional way to turn an AC voltage into DC is to use a bridge rectifier for rectification and then a capacitor for filtering. The combination of bridge rectifier and filter capacitor forms the standard interface circuit that can be used for energy harvesting, i.e. SEH. Ottman et al. discussed the optimization to the SEH technique [3], [4]. Yet, using SEH cannot ensure the energy is always flowing from mechanical part to electrical part. During a certain interval in every cycle, energy returns from electrical part to mechanical part. It was called *energy return phenomenon* [9].

The SSHI treatment overcomes this problem by adding an inductive switching path to the SEH circuit. This path can be connected in parallel to the bridge rectifier to form a p-SSHI circuit (Fig. 2(a)), or in series to form a s-SSHI circuit (Fig. 2(b)) [6]. Regardless of p-SSHI or s-SSHI, the inductive switching path switches on to form a series LCR loop once the displacement  $x$  reaches its extreme values, and then switches off after half of a LCR cycle, as shown in Fig. 3(c), so as to allow a natural inversion to  $v_p$ , the voltage across the piezoelectric element. Typical waveforms of p-SSHI or s-SSHI are shown in Fig. 3. Since  $i_{eq}$  is proportional to  $\dot{x}$ , these switching actions make  $v_p$  in phase with  $i_{eq}$ .

To implement SSHI, additional units for displacement peak detection and switch control are required. In most of the researches, these functions were realized with displacement sensors and digital controllers [5]–[7], all of which need external power to run. In [8], all these units were replaced by a self-powered SSHI circuitry. It makes use of the piezoelectric element itself as a displacement sensor and generates switching commands with transistors accordingly. The introduction of this self-powered SSHI<sup>1</sup> allows the SSHI technique to be more compact and independent of external power, therefore opens a more promising future to the SSHI technique.

<sup>1</sup>More exactly speaking, it is a self-powered series-SSHI device.

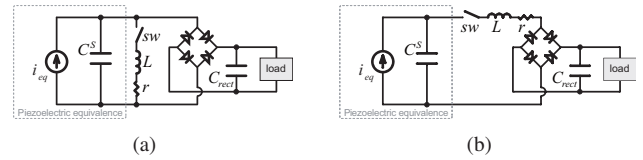


Fig. 2. Equivalent circuits of two SSHI treatments. (a) p-SSHI. (b) s-SSHI.

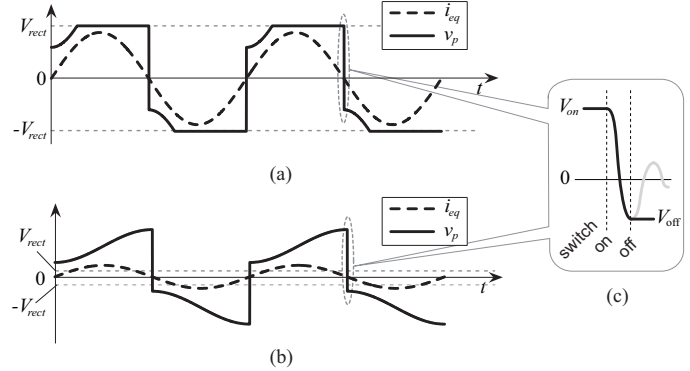


Fig. 3. Typical waveforms of two SSHI treatments. (a) p-SSHI. (b) s-SSHI. (c) Inversion of  $v_p$  at the instant of extreme displacements.

Yet, some concerns about the self-powered SSHI are still unanswered in [8], e.g., is there any constraint or applicable range for this treatment? Does the self-powered SSHI always outperform SEH? In the following sections, the self-powered SSHI treatment will be further investigated concerning these questions.

### III. CIRCUITRY

The essence of the self-powered SSHI technique proposed in [8] is the electronic breaker, which can automatically perform switching action without providing external power when the potential difference across the switch reaches its maximum. Since one breaker can only allow current flow in one direction, replacing the switch  $sw$  in Fig. 2 with two of such breakers (one as maximum breaker and the other, which was inversely connected, as minimum breaker), the self-powered SSHI can be achieved. The breaker consists of three parts: envelope detector, comparator, and switch. In their design, the envelope detectors are in series with the clamped capacitance  $C_p$  and inductor  $L$ . The detected voltage in fact is not  $v_p$ , but the voltage sum of  $v_p$  and the voltage across  $L$ . Even  $L$  is connected to  $C_p$  for a very short interval in every cycle, the hard switching-off action introduces high frequency components to  $L$ . The local maxima or minima produced by these high frequency components may induce misjudgement to the other breaker. Therefore, both envelope detector and comparator parts should be carefully isolated from the switching path. For the breaker introduced in [8], two resistors were connected for isolation purpose. Yet, the principle and design guideline were not clearly addressed.

Taking these envelope detection and isolation issues into consideration, in our self-powered design, we use a complementary transistors topology to achieve both direct envelope

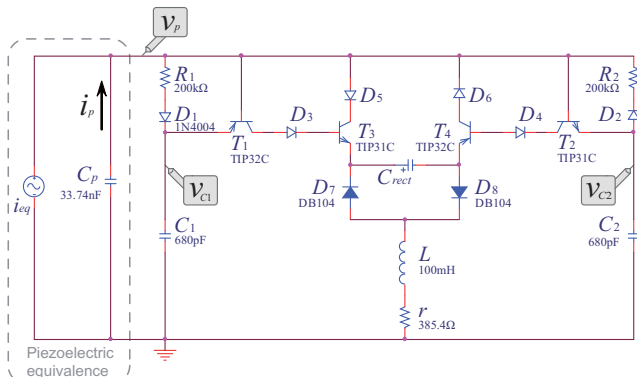


Fig. 4. Modified self-powered SSHI circuitry.

detection for  $v_p$  and reduction on the interference among different parts in the breakers. With this reformation, all isolating resistors, which are bound to consume some energy, can be removed. The modified circuitry is shown in Fig. 4. We obtain its waveforms (Fig. 5), as well as the zoom in view to one of the processes of switching on maximum (Fig. 6) with PSpice simulation. The part values and models given in Fig. 4 are also corresponding to those in the experimental circuit introduced in Section V<sup>2</sup>.

For switching on maximum,  $R_1$ ,  $D_1$  and  $C_1$  form an envelope detector.  $T_1$  and  $T_3$  are cut off most of the time in a cycle. When  $v_p$  reaches its maximum  $V_{max}$ , the voltage across  $C_1$  is  $V_{max} - V_D$ , where  $V_D$  denotes the forward voltage drop of a diode. Then  $v_p$  begins to drop. When the decrease reaches  $V_D + V_{BE}$ , i.e.  $v_p = V_1$  ( $t_1$  instant in Fig. 6),  $T_1$  conducts<sup>3</sup>.  $C_1$  begins to discharge through  $T_{1(ec)}$ ,  $D_3$ ,  $T_{3(be)}$ ,  $C_{rect}$ ,  $D_8$ ,  $L$  and  $r$ , consequently makes  $T_3$  conduct<sup>4</sup>. The conduction of  $T_3$  switches on the inductive path that consists of  $D_5$ ,  $T_{3(ce)}$ ,  $C_{rect}$ ,  $D_8$ ,  $L$  and  $r$ , producing a shortcut to the charge in  $C_p$  and  $C_2$  (through  $D_2$ ,  $R_2$ ). For  $C_p$ , it starts a quick discharge from the voltage of  $V_1$  through the LCR loop, until  $v_p$  reaches its local minimum  $V_2$  ( $t_2$  instant in Fig. 6). The current through  $L$  now tends to reverse its flowing direction, but the  $T_{3(ce)}$  path is immediately blocked by reverse  $D_5$ . Yet, the path consists of  $D_7$ ,  $C_{rect}$ ,  $T_{4(ec)}$ ,  $D_6$  is still available. Because even  $T_4$  is cut off, there is a small parasitic capacitance across its emitter and collector, which is uncharged. The current will stop flowing until  $T_4$ 's emitter-collector capacitance  $C_{CE}$  is charged, at which very instant  $v_p$  becomes  $V_3$  ( $t_3$  instant in Fig. 6). This local minimum of  $v_p$ , i.e.  $V_2$ , may induce misjudgement for the minimum breaker. So  $R_2$  is necessary for making sure that  $C_2$ , which is used for minimum detection, discharges slower than  $C_p$ , so as to skip over this local minimum. After  $t_3$ , both  $T_3$  and  $T_4$  are cut off; however,  $C_2$  still has not finished its discharging, the rest of the charge in  $C_2$  will flow into  $C_p$  and  $C_1$  until they are the same in voltage. This charge neutralization again

<sup>2</sup>Because of the difference between the simulation models and real parts, in simulation, to properly start up the switching processes,  $C_1$  and  $C_2$  are set to be 2nF.

<sup>3</sup> $V_{BE}$  denotes the transistor base-emitter threshold voltage.

<sup>4</sup> $r$  is the equivalent series resistance of  $L$ .

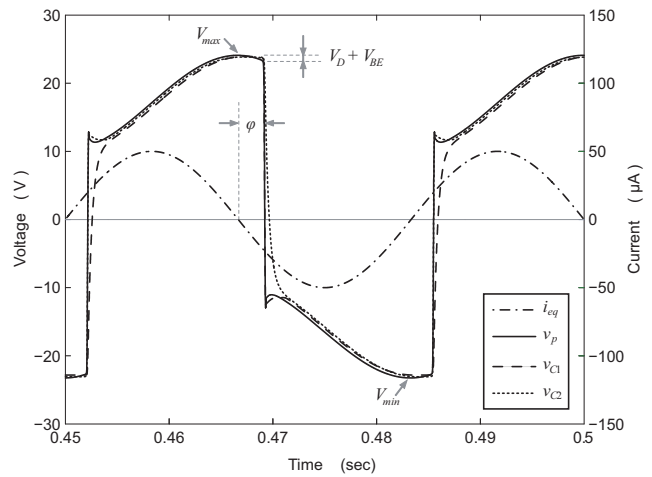


Fig. 5. Simulation waveforms in modified self-powered SSHI.

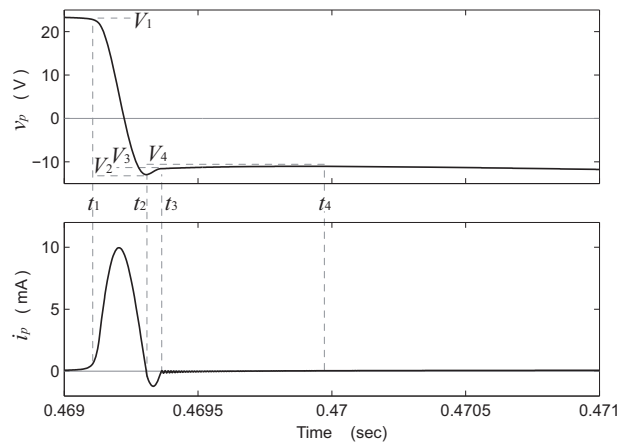


Fig. 6.  $v_p$  and  $i_p$  waveforms in the process of switching on maximum.

increases  $v_p$  a little bit to  $V_4$  ( $t_4$  instant in Fig. 6) before  $v_p$  enters the following half cycle of minimum detection.

Switching on minimum makes use of the counterparts in the circuitry, and its principle is similar to that of switching on maximum.

#### IV. ANALYSIS

In Section III, the working principle of the modified self-powered SSHI circuitry has been introduced. Based on this, detailed and quantitative analysis is provided in this Section.

##### A. Open Circuit Voltage

Regardless of switching on maximum or minimum, the current through  $C_{rect}$  is always positive, so  $C_{rect}$  acts as energy storage. On the contrary, the average power to  $C_1$  and  $C_2$  is zero. They never sustain energy in themselves, so it is unsuitable to regard them as energy storage as did in [8]. Rather,  $C_1$  and  $C_2$  can be equivalently regarded as two capacitors connected in parallel to  $C_p$ . This approximation is validated from Fig. 5, since both  $v_{C1}$  and  $v_{C2}$  are very close to  $v_p$ . Given the harmonic displacement excitation as

$$x(t) = X \sin(\omega t), \quad (2)$$

where  $X$  is the amplitude of maximum displacement,  $\omega$  is the vibration circular frequency. With (1), the equivalent current source should be

$$i_{eq}(t) = \alpha_e X \omega \cos(\omega t). \quad (3)$$

With the parallel connections of  $C_p$ ,  $C_1$  and  $C_2$ , at open circuit condition,  $v_p$  becomes

$$v_{p,oc}(t) = V_{OC} \sin(\omega t), \quad (4)$$

where

$$V_{OC} = \frac{\alpha_e X}{C_p + 2C_{ed}} \quad (5)$$

is the open circuit voltage, representing the amplitude of  $v_{p,oc}$ . Because the capacitances of envelope detecting capacitors  $C_1$  and  $C_2$  are selected to be the same, they are denoted as  $C_{ed}$  in (5). Without shunt circuit connected, the open circuit voltage of the original piezoelectric element is

$$V_{OC,org} = \frac{\alpha_e X}{C_p}. \quad (6)$$

Therefore, (5) implies that the open circuit voltage in self-powered SSHI will be slightly reduced under the same excitation.

In addition, to effectively drive the switches, there is a constraint for  $V_{OC}$ , which is set by the forward voltage gaps of diodes and transistors in the circuit. To figure out the constraint, suppose no any switching action is performed before the connection of the circuit. Once it is connected, the first switching action may start after  $v_p$  attains, for example, its maximum, i.e.  $V_{OC}$ , and then drop to  $V_{OC} - V_D - V_{BE}$ . At this time,  $T_1$  will conduct only when  $v_{C1}$  is larger than the voltage gap produced by  $T_{1(ec)}$ ,  $D_3$ ,  $T_{3(be)}$ ,  $C_{rect}$ ,  $D_8$  in series; and  $T_3$  will conduct only when  $v_p$  is larger than the voltage gap produced by  $D_5$ ,  $T_{3(ce)}$ ,  $C_{rect}$ ,  $D_8$  in series. Both yield the same constraint for  $V_{OC}$  as

$$V_{OC} > V_{CE(sat)} + 3V_D + V_{BE} + V_{DC}, \quad (7)$$

where  $V_{CE(sat)}$  is the collector-emitter saturation voltage of corresponding transistors,  $V_{DC}$  is the rectified voltage, i.e., the voltage across  $C_{rect}$ .

On the other hand, given a  $V_{OC}$  satisfying (7), we can obtain the maximum attainable  $V_{DC}$  in energy harvesting from (7), as follows:

$$V_{DC,max} = V_{OC} - V_{CE(sat)} - 3V_D - V_{BE}. \quad (8)$$

### B. Switching Phase Lag

From the principle of SSHI [6], the switching actions should be taken at the right instants when  $v_p$  attains its extreme values, i.e.  $V_{max}$  or  $V_{min}$  in Fig. 5. In self-powered SSHI, however, to switch at the very instants is impossible; due to the voltage gaps of diode and transistor in envelope detector and comparator, there is always a phase lag between the instants of switching action start and maximum displacement (also  $i_{eq} = 0$ ). The phase lag was defined as  $\varphi$  and regarded as

constant in [8]. Nevertheless,  $\varphi$  in fact changes with  $V_{OC}$ , with the relation of

$$\varphi = \cos^{-1} \left( 1 - \frac{V_D + V_{BE}}{V_{OC}} \right). \quad (9)$$

Considering the constraint on  $V_{OC}$  given in (7), the range of  $\varphi$  can be obtained as

$$0 < \varphi < \cos^{-1} \left[ \frac{V_{CE(sat)} + 2V_D + V_{DC}}{V_{CE(sat)} + 3V_D + V_{BE} + V_{DC}} \right]. \quad (10)$$

The lower limit corresponds to infinite  $V_{OC}$ ; the upper one corresponds to minimum harvestable  $V_{OC}$ .

### C. Voltage Inversion Factor

The voltage inversion factor  $\gamma$  is an important parameter in SSHI. It makes use of the natural oscillation of a LCR circuit, so as to perform a quick inversion for  $v_p$  at right instants. During an inversion, the switch is first closed to enable a LCR loop, and then naturally blocked again after half of a LCR cycle, i.e.

$$\tau = \pi \sqrt{LC}, \quad (11)$$

so that the voltage across the capacitance changes at a short interval  $\tau$ , which is much smaller than the mechanical cycle, with the factor of

$$\gamma = -e^{-\pi/(2Q)}. \quad (12)$$

This definition on voltage inversion factor [9] includes the sign information, therefore is more general, compared to that given in [8].

As described in Section III and illustrated in Fig. 6, the switching process in self-powered SSHI is more complex than that in ordinary SSHI. It might go through two inversion steps and one charge neutralization before the voltage changes from  $V_1$  to  $V_4$ . Among these three steps, there are two intermediate values, which were denominated as  $V_2$  and  $V_3$  in Fig. 6. Taking switching on maximum as example, if  $V_1 > V_{ref1}$ , where

$$V_{ref1} = V_{CE(sat)} + 2V_D + V_{DC} \quad (13)$$

is the first reference voltage gap,  $v_p$  will experience the first inversion. For the first inversion, i.e. from  $V_1$  to  $V_2$ ,  $C_p + C_1$ ,  $L$  and  $r$  form the LCR loop for discharging, with the quality factor of

$$Q_1 = \frac{1}{r} \sqrt{\frac{L}{C_p + C_{ed}}}. \quad (14)$$

The relation between  $V_2$  and  $V_1$  can be obtained as

$$V_2 - V_{ref1} = \gamma_1 (V_1 - V_{ref1}), \quad (15)$$

where

$$\gamma_1 = \begin{cases} -e^{-\pi/(2Q_1)}, & V_1 > V_{ref1}; \\ 1, & \text{others} \end{cases} \quad (16)$$

is the inversion factor for the  $C_p + C_1$ ,  $L$  and  $r$  loop, whose quality factor is  $Q_1$ .

After the first inversion, if  $V_2 < V_{ref2}$ , where

$$V_{ref2} = -2V_D - V_{DC} \quad (17)$$



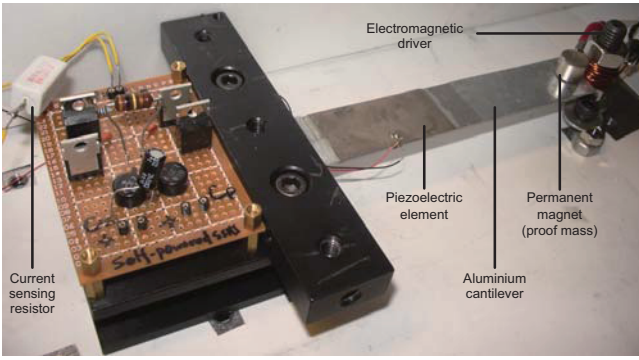


Fig. 7. Experimental setup.

is the second reference voltage gap,  $v_p$  will experience one more inversion. For the second inversion, i.e. from  $V_2$  to  $V_3$ ,  $C_p$  in series with  $C_{CE}$ ,  $L$  and  $r$  form the LCR loop for discharging, with the quality factor of

$$Q_2 = \frac{1}{r} \sqrt{\frac{L(C_p + C_{CE})}{C_p C_{CE}}}. \quad (18)$$

The relation between  $V_3$  and  $V_2$  can be obtained as

$$V_3 - \frac{C_p}{C_{CE}}(V_2 - V_3) - V_{ref2} = \gamma_2(V_2 - V_{ref2}), \quad (19)$$

where

$$\gamma_2 = \begin{cases} -e^{-\pi/(2Q_2)}, & V_{1.5} < V_{ref2}; \\ 1, & \text{others} \end{cases} \quad (20)$$

is the inversion factor for the corresponding LCR loop, whose quality factor is  $Q_2$ .

The charge neutralization follows the second inversion. Since the resistor  $R_2$  is used for slowing down the discharging process of  $C_2$ , roughly speaking, the time constant of  $R_2 C_2$  should be larger than  $\tau$ , which was given in (11). So we can simply assume that the discharge of  $C_2$  starts after the two inversion of  $v_p$ . In the charge neutralization, the total charge in  $C_p$ ,  $C_1$  and  $C_2$  is unchanged. Considering their original voltage,  $V_4$  is related to  $V_1$ ,  $V_2$  and  $V_3$  with the following equation

$$(2C_{ed} + C_p)V_4 = C_{ed}(V_1 + V_2) + C_p V_3. \quad (21)$$

One more relation links  $V_1$ ,  $V_4$  and the open circuit voltage  $V_{OC}$ , i.e.

$$V_1 + V_4 = 2V_{OC} \cos \varphi. \quad (22)$$

$V_1 \sim V_4$  can be expressed in terms of  $V_{OC}$  and  $V_{DC}$  by solving the linear equations of (15), (19), (21) and (22).

Because of the complementary topology, for switching on minimum, the four corresponding voltages are  $-V_1$ ,  $-V_2$ ,  $-V_3$ , and  $-V_4$ , respectively.

#### D. Harvesting Power

Based on the above analyses about the influences of self-powered implementation to the open circuit voltage, switching

TABLE I  
CIRCUIT PARAMETERS.

Name	Symbol	Value
Diode forward voltage drop	$V_D$	0.5 V
Transistor base-emitter on voltage	$V_{BE}$	0.5 V
Transistor collector-emitter saturation voltage	$V_{CE(sat)}$	1.2 V [10]
Transistor emitter-collector capacitance	$C_{CE}$	150 pF [10]
Voltage inversion factor	$\gamma_1$	-0.52

phase lag and voltage inversion factor in SSHI, the analysis on harvesting power can be carried out.

In each vibration cycle, the harvested energy of the self-powered SSHI is

$$E_{SP-SSHI} = 2V_{DC}[C_p(V_1 + V_3 - 2V_2) + C_{ed}(V_1 - V_2)]. \quad (23)$$

Multiplying  $E_{SP-SSHI}$  by the vibration frequency yields the harvesting power of the self-powered SSHI, as

$$P_{SP-SSHI} = f_0 E_{SP-SSHI}, \quad (24)$$

where  $f_0 = \omega/(2\pi)$  is the vibration frequency.

Besides, for SEH and ideal SSHI<sup>5</sup>, the harvesting powers are [6], [8]

$$P_{SEH} = 4f_0 C_p V_{DC}(V_{OC,org} - V_{DC} - 2V_D), \quad (25)$$

$$P_{SSHI} = 4f_0 C_p V_{DC}(V_{OC,org} - V_{DC} - 2V_D) \frac{1-\gamma}{1+\gamma}, \quad (26)$$

respectively.

#### V. EXPERIMENTS

Experiments are performed in order to evaluate the performance of practical self-powered SSHI. The experimental setup is shown in Fig. 7. It is built up with a piezoelectric cantilever and the modified self-power SSHI interface circuitry.

The main mechanical structure is an aluminium cantilever whose fixed end is fixed on the vibration-free table and the free end is driven by an electromagnetic driver. A piezoceramic patch of 49mm × 24mm × 0.508mm (T120-A4E-602, Piezo System, Inc.) is bonded near the fixed end where the largest strain happens along the cantilever. A permanent magnet is attached at the free end of the cantilever, so as to achieve the coupling with the electromagnetic driver; and it also acts as a proof mass to lower the vibration frequency and increase the displacement of the free end. A function generator (33120A, Agilent Co.), following by a power amplifier (2706, B&K Co.), provides a 30Hz sinusoidal excitation to the electromagnetic coil. To perform constant displacement excitation, an inductive displacement sensor (JCW-24SR, CNHF Co.), which is not shown in Fig. 7, is used to sense the displacement of the cantilever for adjustment under different situations.

For the circuitry, component models and values are the same as those shown in Fig. 4. Other circuit parameters are given in Table I.

<sup>5</sup>In ideal SSHI, the sensing and switching control units do not bring any  $V_{OC}$  influence, switching phase delay, and voltage gap to the circuit. However, the voltage gap of bridge rectifier is considered nonzero.

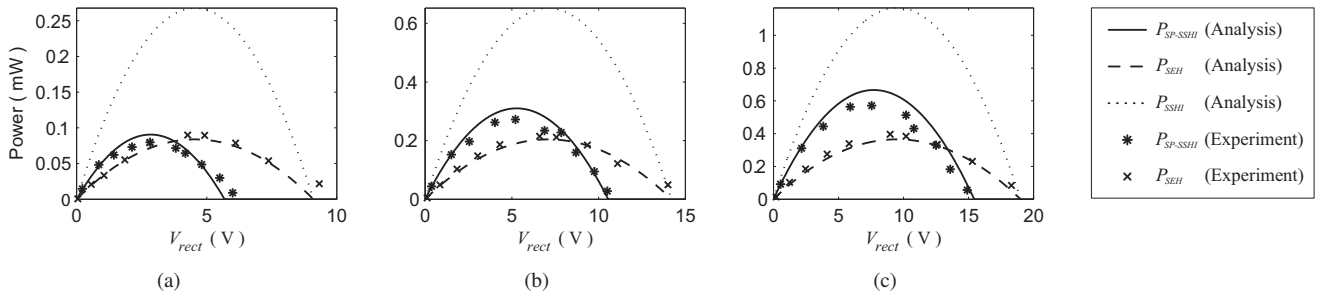


Fig. 8. Harvesting power under different excitation levels. (a)  $V_{OC,org} = 10.1$  V. (b)  $V_{OC,org} = 15.2$  V. (c)  $V_{OC,org} = 20$  V.

TABLE II  
MEASURED CHANGES ON OPEN CIRCUIT VOLTAGE.

$V_{OC,org}$ (V)	$V_{OC}$ (V)	$V_{OC}/V_{OC,org}$
10.1	9.7	0.9604
15.2	14.6	0.9605
20	9.2	0.9600

In experiments, firstly, the changes on open circuit voltage before and after the connection of self-powered SSHI circuit are checked under three excitation levels. As shown in Table II, the ratios of  $V_{OC}/V_{OC,org}$  in these three situations agree with the ratio of  $C_p/(C_p + 2C_{ed}) = 0.9613$  in our experiment, which verified the analysis on open circuit voltage in Section IV.

Also under those three excitation levels, the harvesting power is measured as function of rectified voltage  $V_{DC}$ . Resistors with different resistance values are connected as loads one by one. With the corresponding measured DC voltage across each resistor, the harvesting power under different  $V_{DC}$  can be obtained. The experimental results of  $P_{SP-SSHI}$  and  $P_{SEH}$  under three excitation levels, together with the analyzed  $P_{SP-SSHI}$ ,  $P_{SEH}$  and  $P_{SSHI}$  are given in Fig. 8 for comparison.

From the three sub-figures in Fig. 8, both analytical and experimental results show good agreement with each other. Comparing the self-powered SSHI to the ideal SSHI, the higher the excitation level, the closer between  $P_{SP-SSHI}$  and  $P_{SSHI}$ . On the other hand, comparing the self-powered SSHI to SEH, the maximum harvesting power in self-powered SSHI is larger than that in SEH only when the excitation level is high enough. Therefore, rather than claiming that self-powered SSHI always outperforms SEH, we should note that there should be a critical excitation level, below which this claim is unconvincing.

## VI. CONCLUSION

The introduction of the self-powered version of synchronized switching harvesting on inductor (SSHI) treatment did open a promising territory for piezoelectric energy harvesting with switching technique. Nevertheless, many issues still lie in further improvements on both circuitry and precise modeling. We proposed a modified circuitry for self-powered SSHI. Compared to the circuitry proposed in [8], the modified circuitry not only minimize the interference among different units

in the circuit, so as can enhance the switching performance; but also result in the removal of some resistive components, so as can further diminish the energy dissipation within the switching processes. Improved analysis was carried out considering three aspects of open circuit voltage, switching phase lag, and voltage inversion factor. Unlike the ideal SSHI, which always has better harvesting capability than SEH, it was found from both analyses and experiments that, for self-powered SSHI, only when the excitation level is high enough, it can outperform SEH. Moreover, the higher the excitation level, the more significant the enhancement on harvesting power; therefore, the more beneficial to replace the standard interface with such self-powered switching interface for piezoelectric energy harvesting.

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